

# Magnetic excitations in the spin-5/2 antiferromagnetic trimer substance $\text{SrMn}_3\text{P}_4\text{O}_{14}$

Masashi Hase<sup>1,\*</sup>, Masaaki Matsuda<sup>2,3</sup>, Koji Kaneko<sup>2</sup>, Naoto Metoki<sup>2</sup>, Kazuhisa Kakurai<sup>2</sup>,

Tao Yang<sup>4</sup>, Rihong Cong<sup>4</sup>, Jianhua Lin<sup>4</sup>, Kiyoshi Ozawa<sup>1</sup>, and Hideaki Kitazawa<sup>1</sup>

<sup>1</sup>*National Institute for Materials Science (NIMS),*

*1-2-1 Sengen, Tsukuba, Ibaraki 305-0047, Japan*

<sup>2</sup>*Japan Atomic Energy Agency (JAEA),*

*2-4 Shirakata Shirane, Tokai, Naka, Ibaraki 319-1195, Japan*

<sup>3</sup>*Oak Ridge National Laboratory (ORNL), Oak Ridge, TN 37831, USA*

<sup>4</sup>*College of Chemistry and Molecular Engineering,*

*Peking University, Beijing 100871, People's Republic of China*

(Dated: October 3, 2011)

## Abstract

A quantum-mechanical 1/3 magnetization plateau and magnetic long-range order appear in the large-spin (5/2) substance  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . Magnetization results of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  can be explained by the spin-5/2 isolated antiferromagnetic linear trimer with the intra-trimer interaction ( $J_1$ ) value of 4.0 K. In the present study, to confirm the spin system, we performed inelastic neutron scattering (INS) experiments of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  powders. We observed plural magnetic excitations. The peak positions are 0.46, 0.68, and 1.02 meV. Constant- $Q$ -scan spectra at several  $Q$  values (magnitude of the scattering vector) indicate that the dispersion is weak. The weak dispersion indicates that the excitations are transitions between discrete energy levels. Our INS results are consistent with results expected in the trimer model. We evaluated the  $J_1$  value as 0.29 meV (3.4 K) without considering the other interactions.

PACS numbers: 75.40.Gb, 75.10.Jm, 75.47.Lx, 75.50.Ee

---

\*Electronic address: HASE.Masashi@nims.go.jp

## I. INTRODUCTION

Quantum-mechanical nature is sometimes apparent even in an ordered state of several low-dimensional spin systems formed by small spins. In the triangular antiferromagnet  $\text{CsCuCl}_3$  with spin-1/2, a small jump was observed in the magnetization curve in the magnetic field parallel to the  $c$  axis.[1] This jump was successfully explained as a spin flop process caused by quantum-mechanical effects.[2] In spin-gap systems with spin-1/2 such as the spin-Peierls system in  $\text{CuGeO}_3$  [3, 4] and the two-leg ladder system in  $\text{SrCu}_2\text{O}_3$ , [5, 6] antiferromagnetic long-range order (AF-LRO) appears when small amounts of impurities were doped.[7–11], Nonetheless excitations originating in the singlet-triplet gap in the pure system were observed.[12]

It has not been studied adequately whether the quantum-mechanical nature can remain in an ordered state of spin systems formed by large spins. Quantum-mechanical nature is apparent in the ordered state of the spin-5/2 substance  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . [13] The temperature  $T$  dependence of the magnetic susceptibility indicates occurrence of a magnetic LRO below about 2.6 K. A 1/3 magnetization plateau in magnetization curves was observed at both 1.3 K and 4.2 K (below and above the transition temperature).[13, 14] The plateau was observed in a powder sample. Therefore, the plateau can appear irrespective of the applied magnetic field direction. As a result, the plateau is a quantum-mechanical magnetization plateau generated by an energy gap in the magnetic excitation spectrum.

It is important to determine the origins of the quantum-mechanical nature (magnetization plateau) in  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . We could explain well the magnetic-field  $H$  and  $T$  dependences of the magnetization using the spin-5/2 isolated linear trimer formed by the AF  $J_1$  interaction ( $J_1 = 4.0$  K) depicted as ellipses in Fig.1.[14] Therefore, we consider that the magnetization plateau originates in discrete energy levels of the AF trimer. If the AF trimer model is valid, we can observe magnetic excitations indicating the discrete energy levels even in powder samples. Consequently, we performed inelastic neutron scattering experiments of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  powders.

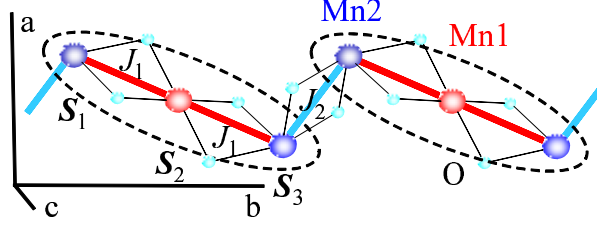


FIG. 1: (Color online) A schematic drawing of positions of  $\text{Mn}^{2+}$  ions ( $3d^5$ ) having localized spin  $5/2$  in  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ .<sup>[13]</sup> Two crystallographically independent  $\text{Mn}^{2+}$  sites (Mn1 and Mn2) exist. Two kinds of short Mn-Mn bonds exist and have Mn-O-Mn paths. The Mn-Mn distances are 3.27 and 3.34 Å at room temperature. The exchange interaction parameters are respectively defined as  $J_1$  and  $J_2$ . The dominant AF  $J_1$  interactions form the spin trimers indicated by the ellipses. The Hamiltonian is expressed as  $\mathcal{H} = J_1(S_1S_2 + S_2S_3)$ . The spin trimer can account for the magnetic-field and temperature dependences of the magnetization when  $J_1 = 4.0$  K. Mn-Mn distances in the other bonds are more than 4.89 Å. These bonds have no Mn-O-Mn paths.

## II. METHODS OF EXPERIMENTS

We synthesized single crystals of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  under hydrothermal conditions at 473 K.<sup>[13]</sup> Each crystal was small. We used pulverized crystals for inelastic neutron scattering (INS) measurements.

We carried out INS measurements on the cold neutron triple-axis spectrometer LTAS installed at JRR-3M in JAEA. The final neutron energy was fixed at 2.6 meV. Higher-order beam contamination was effectively eliminated using a cooled Be filter before the sample. The horizontal collimator sequence was guide-80'-Be-sample-120'-open. This setup yields an energy resolution of 0.1 meV (full width at half maximum, FWHM) at an energy transfer  $\omega = 0$  meV. The resolution was determined from incoherent scattering of the sample. The powder sample of about 9 g was mounted in a  $^4\text{He}$  closed cycle refrigerator.

## III. RESULTS AND DISCUSSION

We performed all the INS measurements above the transition temperature. Circles in Fig. 2 show the  $\omega$  dependence of the INS intensity (constant- $Q$  scan spectra) around 5 K.

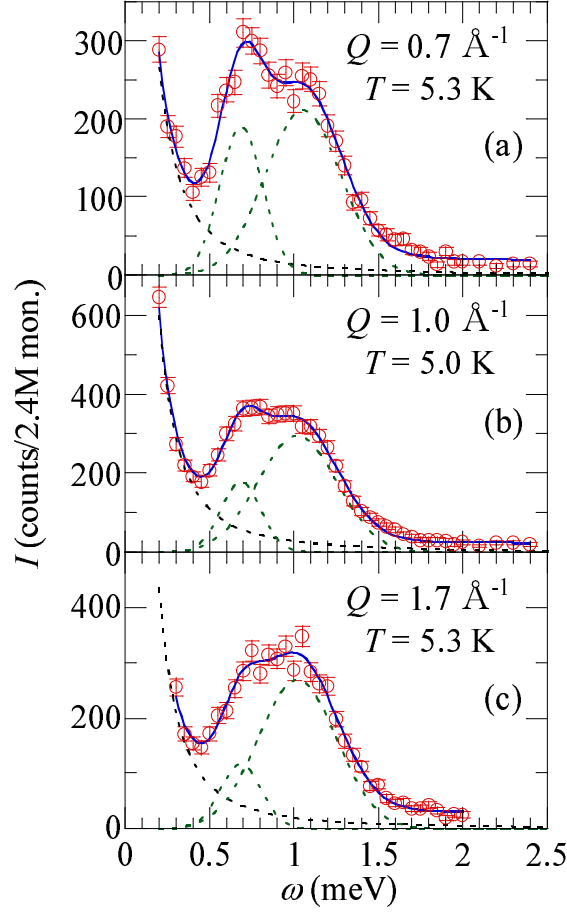


FIG. 2: (Color online) The inelastic neutron scattering intensity *vs.* energy transfer ( $\omega$ ) [constant- $Q$  (magnitude of the scattering vector) scan spectra] of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  around 5 K (circles). The solid line represents the sum of two Gaussians and one Lorentzian (plus constant backgrounds). The dashed line indicates each Gaussian or Lorentzian. The values of the parameters are given in Table I.

The value of  $Q$  is the magnitude of the scattering vector. Excitations are apparent between 0.5 and 1.5 meV. Two kinds of excitations with different peak positions seem to overlap each other. The spectra are almost independent of  $Q$  except for difference in intensities. The weak  $Q$  dependence indicates that the excitations are transitions between discrete energy levels. We consider that the intensities in the vicinity of 0 meV cannot be explained only by incoherent scattering because of  $T$  dependence of spectra as shown later. Low-energy excitations exist.

We compared each spectrum above 0.2 meV in Fig. 2 with a sum of two Gaussians and one Lorentzian (plus constant backgrounds).

$$I(\omega) = \frac{I_0 a_0}{\pi} \frac{1}{\omega^2 + a_0^2} + \sum_i \frac{I_i}{\sqrt{\pi} a_i} \exp\left[-\frac{(\omega - \omega_i)^2}{a_i^2}\right] + I_{\text{BG}}. \quad (1)$$

Here the sum is from  $i = 1$  to 2. The two Gaussians correspond to the excitations between 0.5 and 1.5 meV. The Lorentzian corresponds to the excitations in the vicinity of 0 meV. Each sum of two Gaussians and one Lorentzian (solid line) reproduces well the corresponding spectrum in Fig. 2. Obtained values of the fitting parameters are shown in Table I. The peak position in the spectrum at  $Q = 1.0 \text{ \AA}^{-1}$  is 0.68 meV or 1.02 meV. The peak width (FWHM) is 0.28 meV and 0.55 meV for the 0.68 meV and 1.02 meV excitation, respectively. These widths are larger than the energy resolution of 0.1 meV at  $\omega = 0$  meV, indicating existence of weak dispersion caused by inter-trimer interactions.

Circles in Fig. 3 show constant- $Q$  scan spectra at  $Q = 1.0 \text{ \AA}^{-1}$ . The 0.68 meV and 1.02 meV excitations are also seen at 11.2 K. Intensities around 0.5 meV are larger at 11.2 K than at 5.0 K, suggesting appearance of another transition. Therefore, we compared the spectrum above 0.2 meV at 11.2 K with a sum of three Gaussians and one Lorentzian (plus constant backgrounds) given in Eq. (1) with  $i = 1$  to 3. To reduce variable parameters, we assumed that the peak position (0 meV) of the Lorentzian, and the peak positions (0.68 and 1.02 meV) and widths of the two Gaussians were constant. We used the values obtained at 5.0 K. This assumption is reasonable for transitions between discrete energy levels. The sum of the three Gaussians and one Lorentzian (solid line) reproduces well the spectrum at 11.2 K. The peak position of the third Gaussian is 0.46 meV. The peak width (FWHM) is 0.25 meV and is larger than the energy resolution of 0.1 meV at  $\omega = 0$  meV.

The spectra at 15.6 K and 20.3 K shown respectively in Figs. 3(b) and (c) resemble the spectrum at 11.2 K. Therefore, we compared the spectrum above 0.2 meV at 15.6 K or 20.3 K with a sum of three Gaussians and one Lorentzian (plus constant backgrounds) given in Eq. (1) with  $i = 1$  to 3. In the fitting, we assumed that the peak position of the Lorentzian, and the peak positions and widths of the three Gaussians were constant. The sum (solid line) reproduces well each experimental spectrum. Spectra are featureless above 30 K. We did not compare the spectra with calculated curves. The integrated intensity between 0.5 and 1.5 meV decreases slightly on heating. The Bose factor proportional to phonon intensity at 0.68 meV and 20.3 K, on the other hand, is about 9 times as large as that at 0.68 meV

TABLE I: Values of the integrated intensity ( $I_i$ ) and full width at half maximum (FWHM) of Lorentzian or Gaussian obtained from the fitting of Eq. (1) to experimental constant- $Q$  spectra of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . FWHM is given as  $2a_0$  for Lorentzian or  $2\sqrt{\ln 2}a_i$  for Gaussian. In the fitting to the spectra at  $Q = 1.0 \text{ \AA}^{-1}$ , the value of FWHM of each Gaussian was obtained at 5.0 or 11.2 K and was fixed in the fitting at higher  $T$ . The values in parentheses indicate errors.

$Q$	$T$	0 meV		0.68(1) meV		1.02(2) meV		0.46(5) meV	
		$I_0$	FWHM	$I_1$	FWHM	$I_2$	FWHM	$I_3$	FWHM
$\text{\AA}^{-1}$	K		meV		meV		meV		meV
0.7	5.3	379(58)	0.24(8)	60(9)	0.30(3)	119(11)	0.53(5)		
0.8	5.7	475(136)	0.31(10)	58(8)	0.34(4)	159(21)	0.62(10)		
1.0	5.0	1091(33)	0.16(4)	52(9)	0.28(3)	172(13)	0.55(4)		
1.2	5.4	1123(34)	0.16(4)	32(9)	0.25(5)	190(14)	0.53(4)		
1.5	5.1	968(54)	0.18(3)	31(6)	0.23(3)	177(13)	0.54(3)		
1.7	5.3	795(51)	0.16(4)	36(13)	0.29(6)	157(13)	0.55(6)		
1.0	5.0	1091(33)	0.16(4)	52(9)	0.28(3)	172(13)	0.55(4)		
	11.2	562(296)	0.43(20)	36(6)	0.28	156(12)	0.55	17(12)	0.25(0.14)
	15.6	515(78)	0.53(7)	31(6)	0.28	141(12)	0.55	17(11)	0.25
	20.3	518(179)	0.56(16)	29(6)	0.28	117(11)	0.55	17(10)	0.25

and 5.0 K. Therefore, contribution of phonon is small enough and magnetic excitations are dominant between 0.5 and 1.5 meV.

We examined whether the spin-5/2 AF trimer model with  $J_1 = 4.0$  K can account for the observed excitations. Figures 4(a) depicts a schematic drawing of low-lying energy levels.[14] The following selection rules of transitions are derived theoretically.[15]

$$\Delta S = 0, \pm 1, \quad \Delta M = 0, \pm 1, \quad \text{and} \quad \Delta S_{13} = 0, \pm 1. \quad (2)$$

$S_i (i = 1, 2, 3)$  is the spin operator in the trimer.  $S$  and  $S_{13}$  is defined as  $S_1 + S_2 + S_3$  and  $S_1 + S_3$ , respectively.  $M$  is the  $z$  component of  $S$ . Arrows in Fig. 4(a) indicate allowed transitions from the ground state (GS), first-excited state (1ES) or second-excited states (2ES). In our experimental setup, we can observe transitions with an energy difference  $\Delta\epsilon$  up to 8 when  $J_1 = 4.0$  K. Here,  $\epsilon$  is defined as  $E/J_1$  ( $E$ : eigen energy).

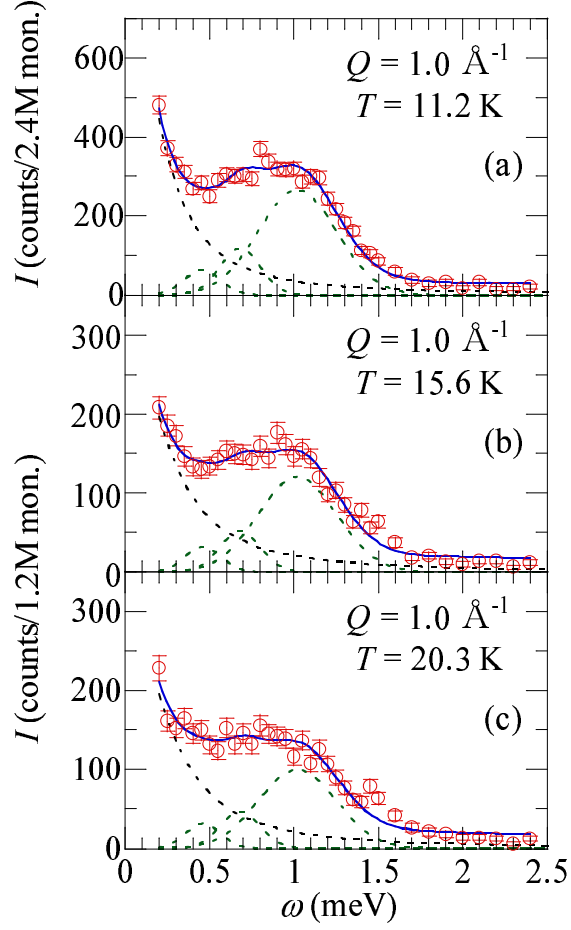


FIG. 3: (Color online) Constant- $Q$  scan spectra of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  at  $Q = 1.0 \text{ \AA}^{-1}$  (circles). The solid line represents the sum of three Gaussians and one Lorentzian (plus constant backgrounds). The dashed line indicates each Gaussian or Lorentzian. The values of the parameters are given in Table I.

Figure 4(b) depicts  $T$  dependence of calculated occupation ratio of the five low-lying energy levels. An inelastic neutron scattering intensity strongly depends on the occupation ratio. From Fig. 4(b), we know that excitations from GS are dominant around 5 K. We considered that the 0.68 and 1.02 meV excitations correspond to transitions from GS to 1ES and 2ES, respectively, indicated by black arrows. The respective energy differences are  $2.5J_1$  and  $3.5J_1$ . The value of  $J_1$  is evaluated as 0.27 meV (3.2 K) or 0.29 meV (3.4 K). These values are slightly smaller than the value determined in the magnetization results ( $J_1 = 4.0$  K). Excitations from 1ES or 2ES are also expected at 11.2 K. We considered that the 0.46

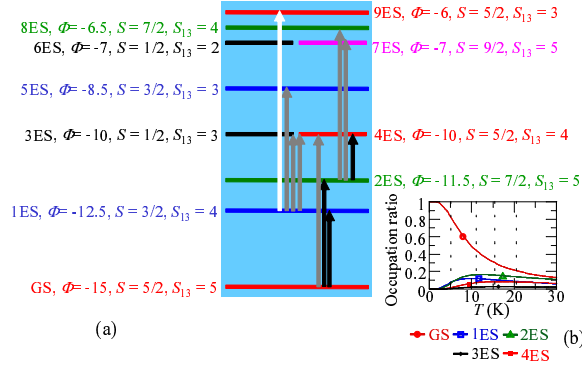


FIG. 4: (Color online) (a) A schematic drawing of low-lying energy levels in the spin-5/2 AF linear trimer [ground state (GS) and excited states (ESs)].[14] The parameters  $\epsilon \equiv E/J_1$  and  $S$  indicate the eigen energy and the total spin, respectively.  $S_{13}$  is defined in the text. To distinguish two degenerate eigen states with  $\epsilon = -10$ , we name the two states 3ES and 4ES. To distinguish two degenerate eigen states with  $\epsilon = -7$ , we name the two states 6ES and 7ES. The arrows indicate allowed transitions from GS, 1ES or 2ES with an energy difference  $\Delta\epsilon$  up to 8. We observed the three black-arrow transitions. The gray-arrow transitions may exist. We could not detect the white-arrow transition. (b) The temperature dependence of calculated occupation ratio of the five low-lying energy levels when  $J_1 = 4.0$  K.

meV excitation corresponds to the transition from 2ES to 4ES indicated by a black arrow. The energy difference is  $1.5J_1$ . The value of  $J_1$  is evaluated as 0.30 meV (3.5 K) and is close to the values evaluated from the other two transitions.

We examined whether we observed all the allowed transitions that are possible in our experimental setup. As was described, we observed the three black-arrow transitions. The gray-arrow transitions may exist. However, we could not prove the existence of the gray-arrow transitions. Energy differences of some gray-arrow transitions ( $2.5J_1$  and  $4J_1$ ) are the same as or close to the energy differences of black-arrow transitions ( $2.5J_1$  and  $3.5J_1$ ). Therefore, we could not extract contribution of the gray-arrow transitions from the experimental results. Energy differences of the other gray-arrow transitions are  $4.5J_1 = 1.3$  meV and  $5J_1 = 1.4$  meV when  $J_1 = 0.29$  meV. Small INS intensities are seen around these energies in Figs. 2 and 3. In our analyses, the small intensities correspond to the tail of the 1.02 meV excitation. However, the peak width is larger in the 1.02 meV excitation (FWHM =



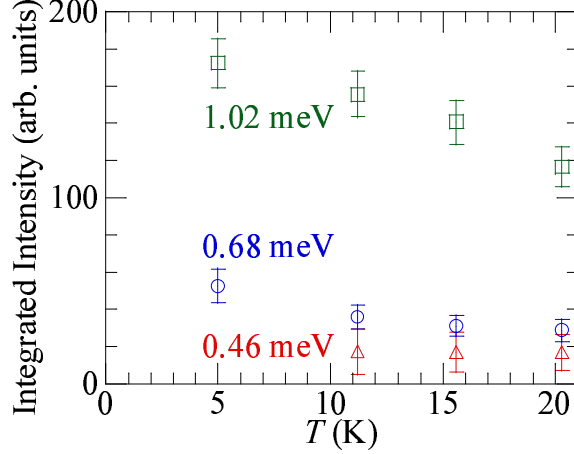


FIG. 5: (Color online) The temperature dependence of the integrated intensity of the 0.46, 0.68, and 1.02 meV excitations in  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  at  $Q = 1.0 \text{ \AA}^{-1}$ .

0.55 meV) than in the 0.46 meV excitation (FWHM = 0.25 meV) or the 0.68 meV excitation (FWHM = 0.28 meV). The gray-arrow transitions may exist in the tail. We could not detect the white-arrow transition. We do not have theoretical INS intensities. Therefore, we could not determine the reason why we could not detect the white-arrow transition. The INS intensity of the white-arrow transition may be very small. Transitions from 3ES or higher excited states must exist at 11.2 K and higher  $T$ . However, we could not prove existence of these transitions because of the same reason for the gray-arrow and white-arrow transitions in Fig. 4(a)

Figure 5 shows the  $T$  dependence of the integrated intensity of the 0.46, 0.68, and 1.02 meV excitations. The integrated intensity of the 0.46 meV excitation is nearly independent of  $T$ . The integrated intensity of the 0.68 or 1.02 meV excitation gradually decreases with increasing  $T$ . As is shown in Fig. 4(a), several transitions are expected to exist. Therefore, the  $T$  dependence of the integrated intensity in Fig. 5 cannot be compared directly with the occupations ratios in Fig. 4(b).

The  $Q$  dependence of the INS intensity in the AF trimer is given in the following formula.[15–17]

$$I(Q) = A_1 f(Q)^2 [1 - \sin(3.27Q)/(3.27Q)] + A_2 f(Q)^2 [1 - \sin(6.54Q)/(6.54Q)]. \quad (3)$$

The values 3.27 and 6.54 indicates the Mn1-Mn2 and Mn2-Mn2 length in the AF trimer,

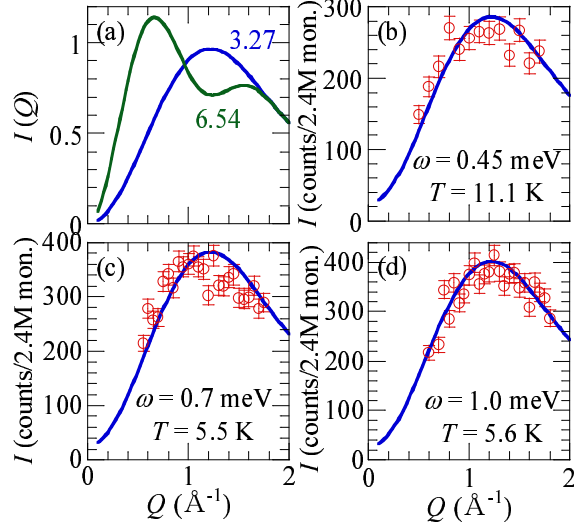


FIG. 6: (Color online) (a) The  $Q$  dependence of the calculated INS intensity in the spin-5/2 AF linear trimer. (b)-(d) The inelastic neutron scattering intensity *vs.*  $Q$  (constant- $\omega$  scan spectra) of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  (circles). The line indicates the first term in Eq. (3).

respectively. The function  $f(Q)$  is the magnetic form factor of  $\text{Mn}^{2+}$  ions.[18] The coefficients  $A_1$  and  $A_2$  depend on two eigen states between which the transition occurs. The coefficients are not derived theoretically. Figure 6(a) represents the two terms in Eq. (3). Circles in Figs. 6(b) - (d) show constant- $\omega$  scan spectra of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . The lines indicate the first term in Eq. (3) plus constant backgrounds. The INS intensity in the vicinity of  $\omega = 2$  meV is small and almost independent of  $\omega$ ,  $Q$ , and  $T$ . Therefore, we used the intensity at  $\omega = 2$  meV for the value of constant backgrounds. The each line is consistent with the corresponding constant- $\omega$  scan spectrum. If  $A_1$  is much larger than  $A_2$  in the observed transitions, this consistency indicates that the AF trimer model can explain the experimental  $I(Q)$ .

We comment on inter-trimer interactions. The dispersion relation of magnetic excitations was calculated in spin dimers with weak inter-dimer interactions using random phase approximation.[19] A similar dispersion relation was inferred in interacting spin tetramers.[20] According to the results, we speculate that the following dispersion relation may be applicable to spin trimers with weak inter-trimer interactions.

$$\omega_{\mathbf{q}=(h,k,l)} = \sqrt{\Delta^2 + \alpha \Delta J(\mathbf{q}) R(T)}. \quad (4)$$

Here,  $\Delta$  is an energy difference between ground and excited states.  $\alpha$  is a coefficient derived

from transition matrix elements. The value of  $\alpha$  is 2 for the spin-1/2 dimer [21] or 5 for the spin-3/2 dimer.[19]  $J(\mathbf{q})$  is a Fourier transform of inter-trimer interactions.  $R(T)$  is a difference in thermal populations of ground and excited states. We consider that the dominant inter-trimer interaction is the  $J_2$  interaction.  $J(\mathbf{q})$  is expressed approximately as  $2J_2\cos(2\pi k)$ . We assume that excitation energies are the same at 0 and 5.0 K. The excitation energy at 0 K at the bottom of the dispersion  $\omega_b$ , where the INS intensity is the strongest, is expressed as follows.

$$\omega_b = \sqrt{\Delta^2 - 2\alpha\Delta J_2}. \quad (5)$$

We assume that only the  $J_2$  interaction is the origin of the difference between the expected excitation energy  $2.5J_1 = 10$  K and experimental excitation energy  $0.68 \text{ meV} = 7.9$  K in the transition between GS and 1ES. Using  $\Delta = 10$  K and  $\omega_b = 7.9$  K, we obtained  $\alpha J_2 = 1.9$  K. If  $\alpha$  is large, a  $J_2$  value can be small enough in comparison with the  $J_1$  value.

We observed magnetic excitations that are consistent with excitations expected in the spin-5/2 AF trimer. Therefore, the discrete energy levels of the AF trimer are the origins of the quantum-mechanical nature (magnetization plateau) in  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . In the strict sense, the energy difference between GS ( $S=5/2$ ) and 2ES ( $S=7/2$ ) generates the magnetization plateau. The magnetization plateau appears even in the ordered state.[13, 14] The property of the cluster (trimer in this case) can remain in the ordered state. The total spin of the ground state of the AF trimer is finite ( $5/2$ ). Therefore, we consider that the magnetic LRO is stabilized by the  $J_1$  and weak three-dimensional inter-trimer interactions. Several cluster substances can maintain their cluster properties in their ordered states. For example, the 1/2 quantum-mechanical magnetization plateau is generated by discrete energy levels of a spin-1/2 tetramer in  $\text{Cu}_2\text{CdB}_2\text{O}_6$ . [22, 23] The plateau remains in the ordered state. With the aid of other researchers, some of the present authors determined the magnetic structure below the transition temperature of  $T_N = 2.2(1)$  K using neutron powder diffraction data.[24] The magnetic structure has a long-range period. We are now considering the origin of the long-range period. We will report details of the magnetic structure in a subsequent paper.

We comment on the INS intensity in the vicinity of 0 meV. Figure 7 shows constant- $Q$  scan spectra at  $Q = 1.0 \text{ \AA}^{-1}$  below 0.5 meV. As was described, the increase of the intensity around 0.4 meV is caused by the 0.46 meV excitation. The intensity between 0.1 and 0.2 meV decreases with increasing  $T$ . This temperature dependence cannot be explained by incoherent scattering or phonon. In addition, we observed diffuse scattering between  $2\theta = 15$

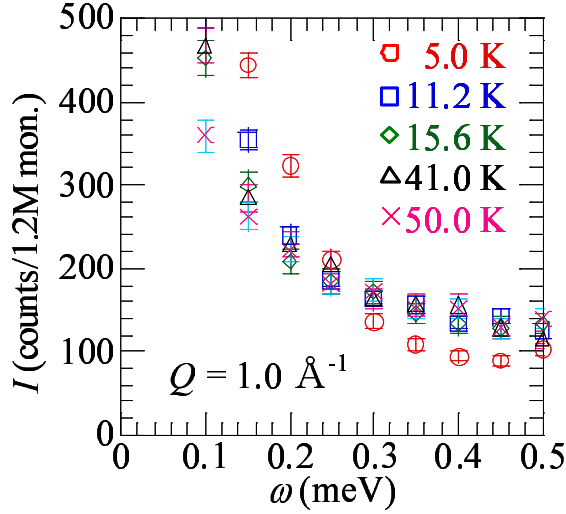


FIG. 7: (Color online) Constant- $Q$  scan spectra of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  at  $Q = 1.0 \text{ \AA}^{-1}$ .

and  $40^\circ$  in neutron powder diffraction patterns (wavelength  $\lambda = 2.458 \text{ \AA}$ ). [24] This  $2\theta$  range corresponds to  $Q = 0.7$  to  $1.7 \text{ \AA}^{-1}$  in the present INS experiments. The shape of the diffuse scattering resembles one-dimensional or two-dimensional Bragg scattering with a cutoff at low  $Q$  and long tail at large  $Q$ . The integrated intensity of the diffuse scattering shows a maximum in the vicinity of  $T_N$ . Several magnetic reflections appear below  $T_N$  between  $2\theta = 15$  and  $40^\circ$ . Therefore, the origin of the diffuse scattering is magnetic. Consequently, magnetic excitations exist in the vicinity of  $0 \text{ meV}$ . The magnetic excitations cannot be explained by transitions between energy levels in the trimer. Spin fluctuation in the ground state generates the magnetic excitations. Therefore, we used the Lorentzian with the  $0 \text{ meV}$  peak in the fitting of the constant- $Q$  scan spectra in Figs. 2 and 3. In future studies, we will perform INS measurements of  $\text{SrMn}_3\text{P}_4\text{O}_{14}$  in the ordered state. Anisotropy of the  $\text{Mn}^{2+}$  spins is small. [14] Therefore, a gap of spin-wave excitations is small. We may observe spin-wave excitations in the vicinity of  $0 \text{ meV}$  in addition to the trimer excitations.

#### IV. SUMMARY

In order to confirm the spin system, we performed inelastic neutron scattering (INS) experiments of powders of the spin-5/2 antiferromagnetic trimer substance  $\text{SrMn}_3\text{P}_4\text{O}_{14}$ . We observed plural magnetic excitations. The peak positions are  $0.46$ ,  $0.68$ , and  $1.02 \text{ meV}$ . The

weak  $Q$  dependence of constant- $Q$ -scan spectra indicates that the excitations are transitions between discrete energy levels. The experimental results are consistent with results expected in the trimer model with the intra-trimer interaction value of 0.29 meV (3.4 K) without considering the other interactions.

## Acknowledgments

We are grateful to T. Masuda for invaluable discussion. The neutron scattering experiments were carried out in the framework of JAEA Users' Program and within the NIMS-RIKEN-JAEA Cooperative Research Program on Quantum Beam Science and Technology. This work was partially supported by grants from NIMS.

- 
- [1] H. Nojiri, Y. Tokunaga, and M. Motokawa, J. Phys. (Paris), **49**, 1459 (1988).
  - [2] T. Nikuni and H. Shiba, J. Phys. Soc. Jpn. **62**, 3268 (1993).
  - [3] M. Hase, I. Terasaki, and K. Uchinokura, Phys. Rev. Lett. **70**, 3651 (1993).
  - [4] M. Hase, I. Terasaki, K. Uchinokura, M. Tokunaga, N. Miura, and H. Obara, Phys. Rev. B **48**, 9616 (1993).
  - [5] E. Dagotto, J. Riera, and D. Scalapino, Phys. Rev. B **45**, R5744 (1992).
  - [6] M. Azuma, Z. Hiroi, M. Takano, K. Ishida, and Y. Kitaoka, Phys. Rev. Lett. **73**, 3463 (1994).
  - [7] M. Hase, I. Terasaki, Y. Sasago, K. Uchinokura, and H. Obara, Phys. Rev. Lett. **71**, 4059 (1993).
  - [8] M. Hase, N. Koide, K. Manabe, Y. Sasago, K. Uchinokura, and A. Sawa, Physica B **215**, 164 (1995).
  - [9] M. Hase, K. Uchinokura, R. J. Birgeneau, K. Hirota, and G. Shirane, J. Phys. Soc. Jpn. **65**, 1392 (1996).
  - [10] M. Azuma, Y. Fujishiro, M. Takano, M. Nohara, and H. Takagi, Phys. Rev. B **55**, R8658 (1997).
  - [11] M. Azuma, M. Takano, and R. S. Eccleston, J. Phys. Soc. Jpn. **67**, 740 (1998).
  - [12] M. C. Martin, M. Hase, K. Hirota, G. Shirane, Y. Sasago, N. Koide, and K. Uchinokura, Phys. Rev. B **56**, 3173 (1997).

- [13] T. Yang, Y. Zhang, S. Yang, G. Li, M. Xiong, F. Liao, and J. Lin, *Inorg. Chem.* **47**, 2562 (2008).
- [14] M. Hase, T. Yang, R. Cong, J. Lin, A. Matsuo, K. Kindo, K. Ozawa, and H. Kitazawa, *Phys. Rev. B* **80**, 054402 (2009).
- [15] A. Furrer and H. U. Güdel, *J. Magn. Magn. Mater.* **14**, 256 (1979).
- [16] M. Matsuda, K. Kakurai, A. A. Belik, M. Azuma, M. Takano, M. Fujita, *Phys. Rev. B* **71**, 144411 (2005).
- [17] A. Podlesnyak, V. Y. Pomjakushin, E. V. Pomjakushina, K. Conder, and A. Furrer, *Phys. Rev. B* **76**, 064420 (2007).
- [18] P. J. Brown, *International Tables for Crystallography Vol. C* (Edited by E. Prince), chapter 4.4.5 (Magnetic form factors), 454.
- [19] B. Leuenberger, A. Stebler, H. U. Güdel, A. Furrer, R. Feile, and J. K. Kjems, *Phys. Rev. B* **30**, 6300 (1984).
- [20] P. S. Häfliger, S. T. Ochsenbein, B. Trusch, H. U. Güdel, and A. Furrer, *J. Phys.: Condens. Matter* **21**, 026019 (2009).
- [21] Y. Sasago, K. Uchinokura, A. Zheludev, and G. Shirane *Phys. Rev. B* **55**, 8357 (1997).
- [22] M. Hase, M. Kohno, H. Kitazawa, O. Suzuki, K. Ozawa, G. Kido, M. Imai, and X. Hu, *Phys. Rev. B* **72**, 172412 (2005).
- [23] M. Hase, A. Dönni, V. Yu. Pomjakushin, L. Keller, F. Gozzo, A. Cervellino, and M. Kohno, *Phys. Rev. B* **80**, 104405 (2009).
- [24] M. Hase *et al.*, unpublished.